

Mechanisms for Generating Power-Law Size Distributions, Part 1

Principles of Complex Systems | @pocsvox
CSYS/MATH 300, Fall, 2017

Prof. Peter Dodds | @peterdodds

Dept. of Mathematics & Statistics | Vermont Complex Systems Center
Vermont Advanced Computing Core | University of Vermont



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Mechanisms, Pt. 1

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Productions



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The First Return Problem
Examples

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Basics
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PLIPLO

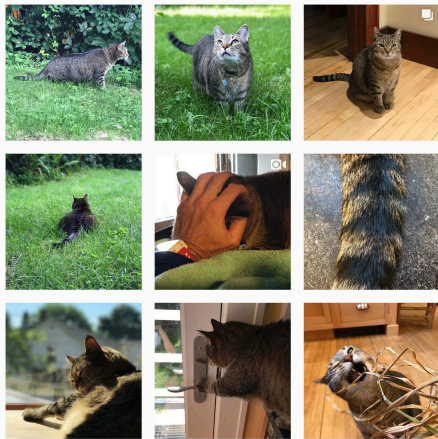
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Special Guest Executive Producer: Pratchett





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 On Instagram at [pratchett_the_cat](https://www.instagram.com/pratchett_the_cat) 



Outline

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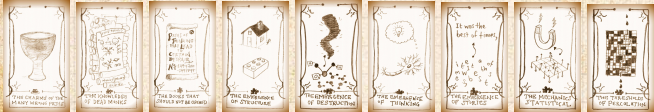
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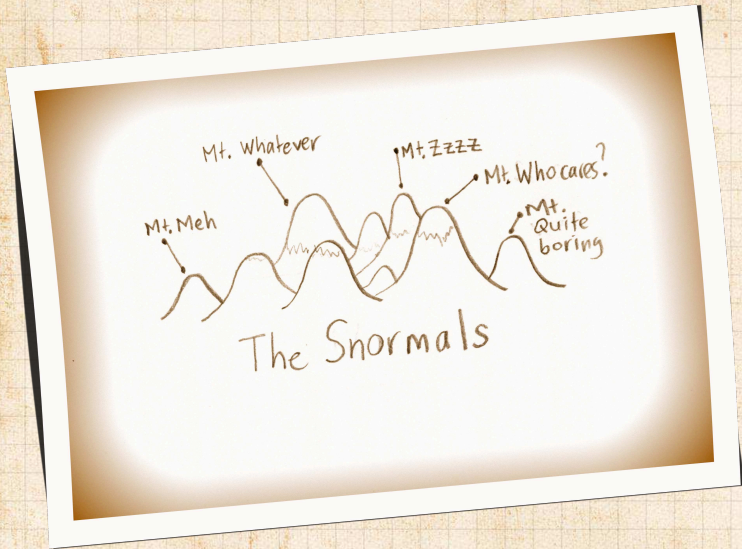
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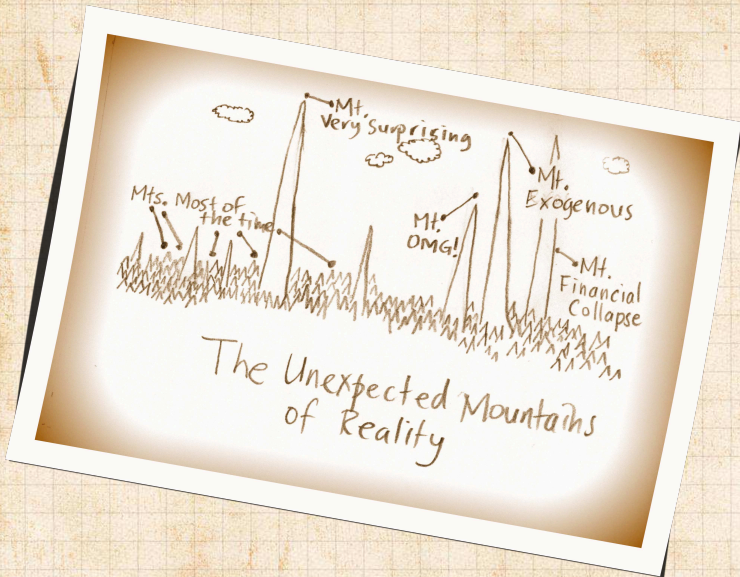
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Mechanisms:

A powerful story in the rise of complexity:

🌀 structure arises out of randomness.

🌀 Exhibit A: Random walks. ↗

The essential random walk:

🌀 One spatial dimension.

🌀 Time and space are discrete.

🌀 Random walker (e.g., a drunk) starts at origin

$$x = 0.$$

🌀 Step at time t is x_t .

$$x_t = \begin{cases} -1 & \text{with probability } 1/2 \\ +1 & \text{with probability } 1/2 \end{cases}$$

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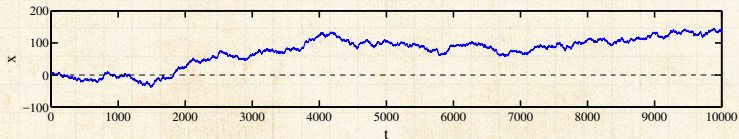
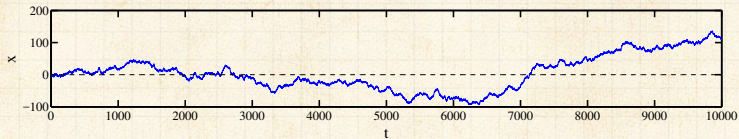
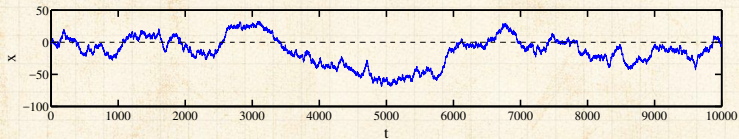
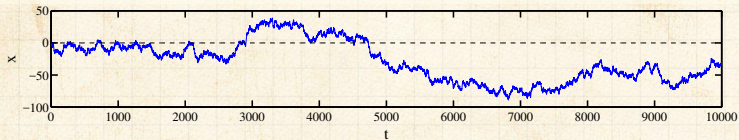
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A few random random walks:

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Random walks:

Displacement after t steps:

$$x_t = \sum_{i=1}^t \epsilon_i$$

Expected displacement:

$$\langle x_t \rangle = \left\langle \sum_{i=1}^t \epsilon_i \right\rangle = \sum_{i=1}^t \langle \epsilon_i \rangle = 0$$

- At any time step, we 'expect' our drunkard to be back at the pub.
- Obviously fails for odd number of steps...
- But as time goes on, the chance of our drunkard lurching back to the pub must diminish, right?

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Variances sum: *

$$\begin{aligned}\text{Var}(x_t) &= \text{Var}\left(\sum_{i=1}^t \epsilon_i\right) \\ &= \sum_{i=1}^t \text{Var}(\epsilon_i) = \sum_{i=1}^t 1 = t\end{aligned}$$

* Sum rule = a good reason for using the variance to measure spread; only works for independent distributions.

So typical displacement from the origin scales as:

$$\sigma = t^{1/2}$$

 A non-trivial scaling law arises out of additive aggregation or accumulation.

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Variations sum: *

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Stock Market randomness:

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Also known as the bean machine ↗, the quincunx
(simulation) ↗, and the Galton box.



Great moments in Televised Random Walks:

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Plinko!  from the Price is Right.



Random walk basics:

Counting random walks:

- Each **specific** random walk of length t appears with a chance $1/2^t$.
- We'll be more interested in how many random walks end up at the same place.
- Define $N(i, j, t)$ as # distinct walks that start at $x = i$ and end at $x = j$ after t time steps.
- Random walk must displace by $+(j - i)$ after t steps.
- Insert question from assignment 3

$$N(i, j, t) = \binom{t}{(t + j - i)/2}$$

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How does $P(x_t)$ behave for large t ?

Take time $t = 2n$ to help ourselves.

$x_{2n} \in \{0, \pm 2, \pm 4, \dots, \pm 2n\}$

x_{2n} is even so set $x_{2n} = 2k$.

Using our expression $N(i, j, t)$ with $i = 0, j = 2k$, and $t = 2n$, we have

$$\Pr(x_{2n} = 2k) \propto \binom{2n}{n+k}$$

For large n , the binomial deliciously approaches the Normal Distribution of Snoredom:

$$\Pr(x_t = x) \simeq \frac{1}{\sqrt{2\pi t}} e^{-\frac{x^2}{2t}}$$

Insert question from assignment 3

The whole is different from the parts. #nutritious

See also: Stable Distributions

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 See also: [Gaussian Distributions](#) 




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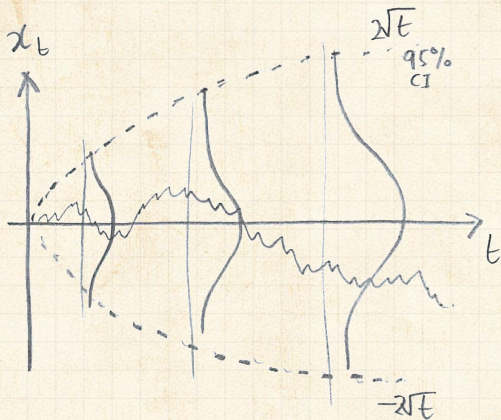
 See also: Stable Distributions 



Universality is also not left-handed:

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
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This is Diffusion : the most essential kind of spreading (more later).



View as Random Additive Growth Mechanism.



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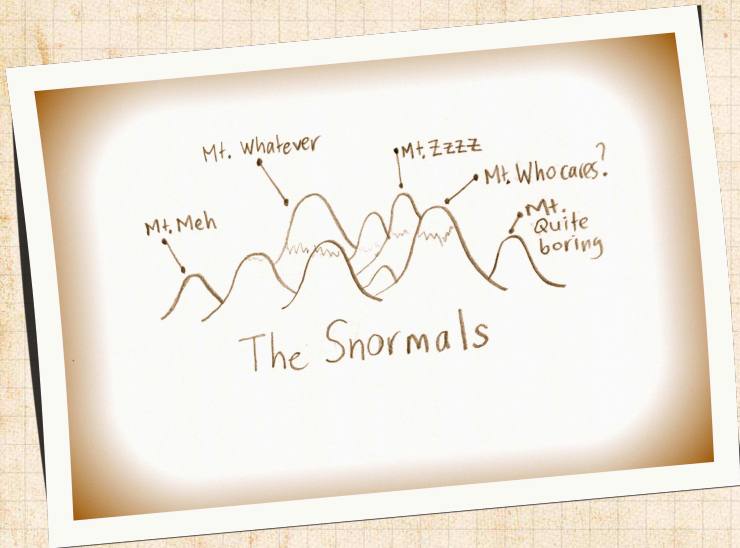
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Random walks are even weirder than you might think...

- ❧ $\xi_{r,t}$ = the probability that by time step t , a random walk has crossed the origin r times.
- ❧ Think of a coin flip game with ten thousand tosses.
- ❧ If you are behind early on, what are the chances you will make a comeback?
- ❧ The most likely number of lead changes is...
- ❧ In fact: $\xi_{0,t} > \xi_{1,t} > \xi_{2,t} > \dots$
- ❧ Even crazier:
The expected time between tied scores = ∞

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See Feller, Intro to Probability Theory, Volume I [3]

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Applied knot theory:



"Designing tie knots by random walks" ↗

Fink and Mao,
Nature, **398**, 31–32, 1999. [4]

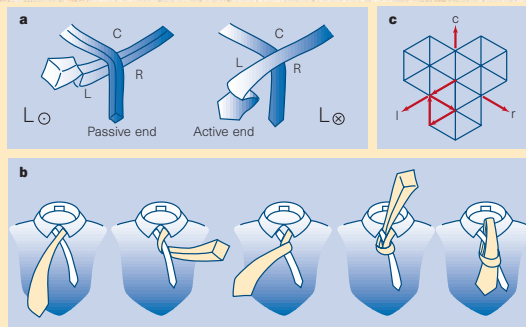


Figure 1 All diagrams are drawn in the frame of reference of the mirror image of the actual tie.
a. The two ways of beginning a knot, L_{\ominus} and L_{\otimes} . For knots beginning with L_{\ominus} , the tie must begin inside-out. **b.** The four-in-hand, denoted by the sequence $L_{\otimes}, R_{\ominus}, L_{\ominus}, C, T$. **c.** A knot may be represented by a persistent random walk on a triangular lattice. The example shown is the four-in-hand, indicated by the walk $\uparrow\uparrow\uparrow\uparrow$.


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Table 1 **Aesthetic tie knots**


h	γ	γ/h	$K(h, \gamma)$	s	b	Name	Sequence
3	1	0.33	1	0	0		$L_{\circ}R_{\circ}C_{\circ}T$
4	1	0.25	1	-1	1	Four-in-hand	$L_{\circ}R_{\circ}L_{\circ}C_{\circ}T$
5	2	0.40	2	-1	0	Pratt knot	$L_{\circ}C_{\circ}R_{\circ}L_{\circ}C_{\circ}T$
6	2	0.33	4	0	0	Half-Windsor	$L_{\circ}R_{\circ}C_{\circ}L_{\circ}R_{\circ}C_{\circ}T$
7	2	0.29	6	-1	1		$L_{\circ}R_{\circ}L_{\circ}C_{\circ}R_{\circ}L_{\circ}C_{\circ}T$
7	3	0.43	4	0	1		$L_{\circ}C_{\circ}R_{\circ}C_{\circ}L_{\circ}R_{\circ}C_{\circ}T$
8	2	0.25	8	0	2		$L_{\circ}R_{\circ}L_{\circ}C_{\circ}R_{\circ}L_{\circ}R_{\circ}C_{\circ}T$
8	3	0.38	12	-1	0	Windsor	$L_{\circ}C_{\circ}R_{\circ}L_{\circ}C_{\circ}R_{\circ}L_{\circ}C_{\circ}T$
9	3	0.33	24	0	0		$L_{\circ}R_{\circ}C_{\circ}L_{\circ}R_{\circ}C_{\circ}L_{\circ}R_{\circ}C_{\circ}T$
9	4	0.44	8	-1	2		$L_{\circ}C_{\circ}R_{\circ}C_{\circ}L_{\circ}C_{\circ}R_{\circ}L_{\circ}C_{\circ}T$


Knots are characterized by half-winding number h , centre number γ , centre fraction γ/h , knots per class $K(h, \gamma)$, symmetry s , balance b , name and sequence.

 h = number of moves

 γ = number of center moves

 $K(h, \gamma) = 2^{\gamma-1} \binom{h-\gamma-2}{\gamma-1}$

 $s = \sum_{i=1}^h x_i$ where $x = -1$ for L and $+1$ for R .

 $b = \frac{1}{2} \sum_{i=2}^{h-1} |\omega_i + \omega_{i-1}|$
where $\omega = \pm 1$
represents winding direction.

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Random walks #crazytownbananapants

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Power-Law
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The problem of first return:

- 1. What is the probability that a random walker in one dimension returns to the origin for the first time after t steps?
- 2. Will our drunkard always return to the origin?
- 3. What about higher dimensions?

Reasons for caring:

1. We will find a power-law size distribution with an interesting exponent.
2. Some physical structures may result from random walks.
3. We'll start to see how different scalings relate to each other.

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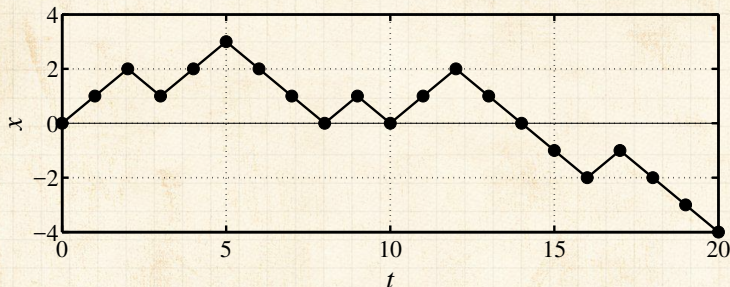
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For random walks in 1-d:



- A **return** to origin can only happen when $t = 2n$.
- In example above, returns occur at $t = 8, 10,$ and 14 .
- Call $P_{fr(2n)}$ the probability of **first return** at $t = 2n$.
- Probability calculation \equiv Counting problem (combinatorics/statistical mechanics).
- Idea: Transform first return problem into an easier return problem.

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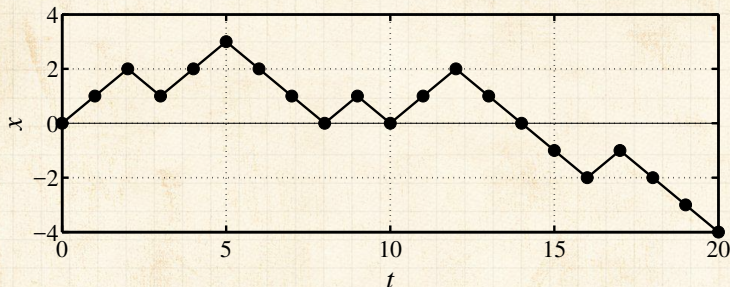
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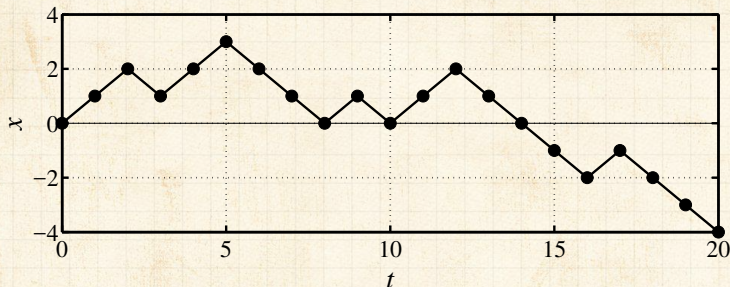
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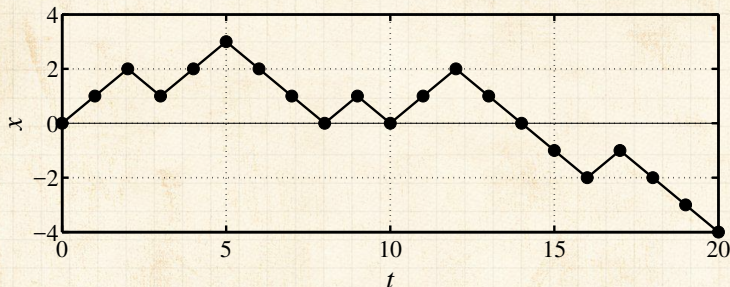
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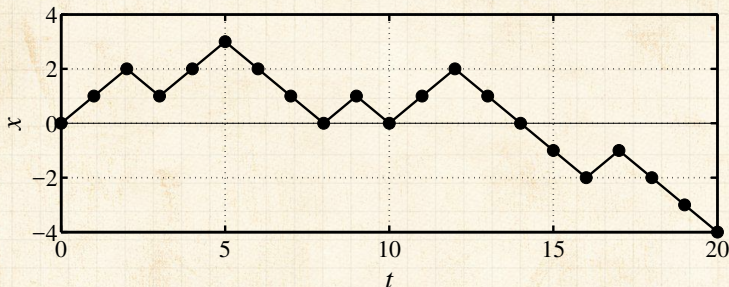
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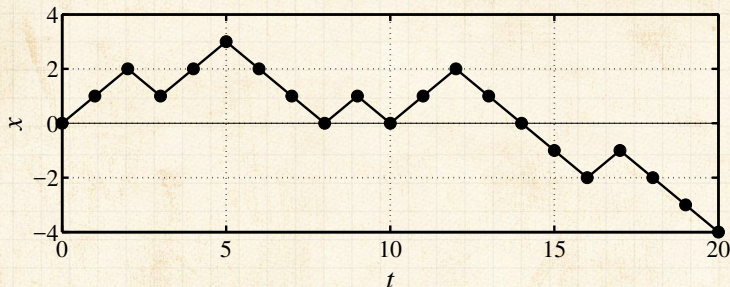
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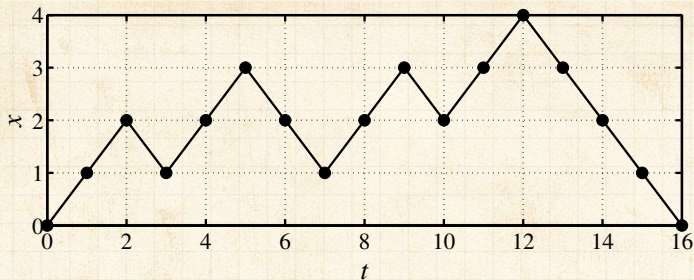
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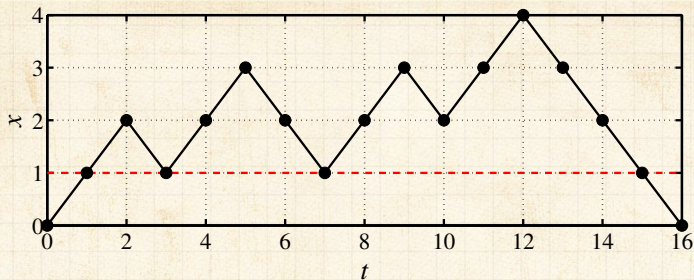
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- Can assume drunkard first lurches to $x = 1$.
- Observe walk first returning at $t = 16$ stays at or above $x = 1$ for $1 \leq t \leq 15$ (dashed red line).
- Now want walks that can return many times to $x = 1$.
- $P_{fr}(2n) = 2 \cdot \frac{1}{2} Pr(x_t \geq 1, 1 \leq t \leq 2n-1, \text{ and } x_1 = x_{2n-1} = 1)$
- The $\frac{1}{2}$ accounts for $x_{2n} = 2$ instead of 0.
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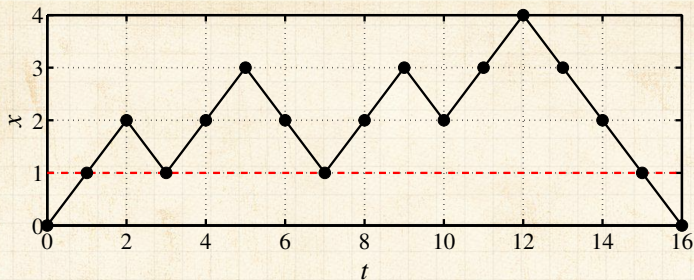
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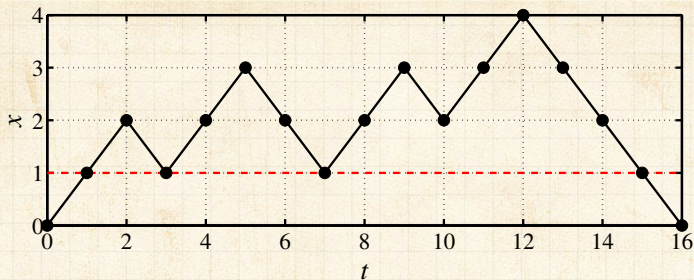


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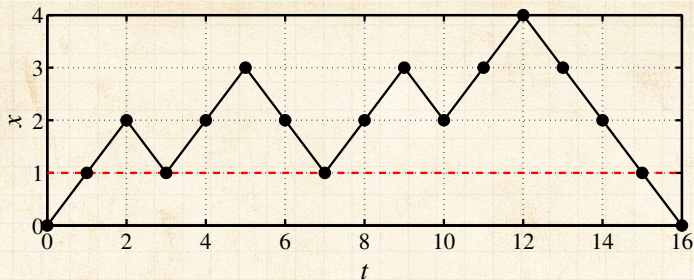




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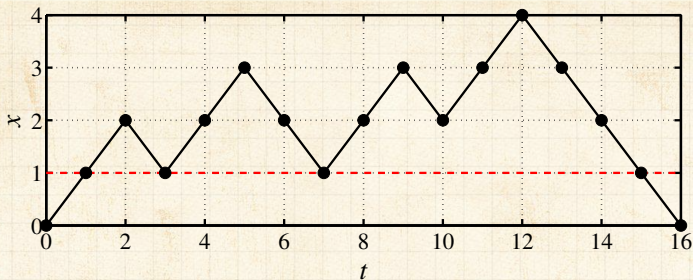




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Counting first returns:

Approach:

- Move to counting numbers of walks.
- Return to probability at end.
- Again, $N(i, j, t)$ is the # of possible walks between $x = i$ and $x = j$ taking t steps.
- Consider all paths starting at $x = 1$ and ending at $x = 1$ after $t = 2n - 2$ steps.
- Idea: If we can compute the number of walks that hit $x = 0$ at least once, then we can subtract this from the total number to find the ones that maintain $x \geq 1$.
- Call walks that drop below $x = 1$ excluded walks.
- We'll use a method of images to identify these excluded walks.

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- Consider **all paths** starting at $x = 1$ and ending at $x = 1$ after $t = 2n - 2$ steps.

- Idea: If we can compute the number of walks that hit $x = 0$ at least once, then we can subtract this from the total number to find the ones that maintain $x \geq 1$.
- Call walks that drop below $x = 1$ **excluded walks**.
- We'll use a method of images to identify these excluded walks.

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Counting first returns:

Approach:

- Move to counting numbers of walks.
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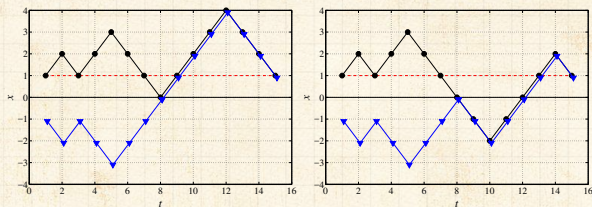
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Examples of excluded walks:



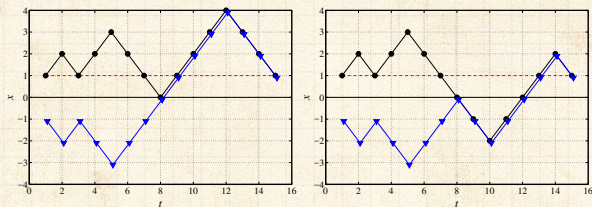
Key observation for excluded walks:

- For any path starting at $x=1$ that hits 0, there is a unique matching path starting at $x=-1$.
- Matching path first mirrors and then tracks after first reaching $x=0$.
- # of t -step paths starting and ending at $x=1$ and hitting $x=0$ at least once

So $N_{\text{first return}}(2n) = N(1, 1, 2n-2) - N(-1, 1, 2n-2)$



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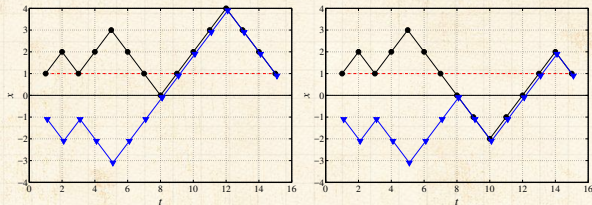
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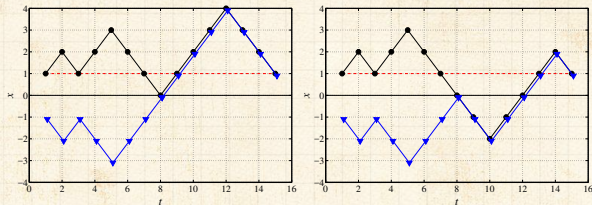
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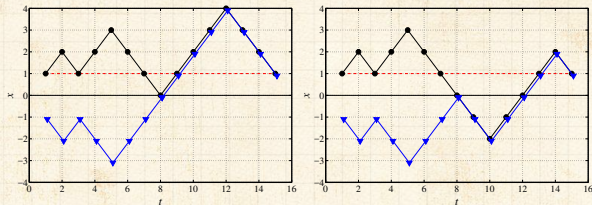
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Probability of first return:

Insert question from assignment 3  :

 Find

$$N_{\text{fr}}(2n) \sim \frac{2^{2n-3/2}}{\sqrt{2\pi n}^{3/2}}$$

 Normalized number of paths gives probability.

 Total number of possible paths = 2^{2n} .



$$P_{\text{fr}}(2n) = \frac{1}{2^{2n}} N_{\text{fr}}(2n)$$

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Higher dimensions!

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Random walks

On finite spaces:

- 1. In any finite homogeneous space, a random walker will visit every site with equal probability
- 2. Call this probability the **Invariant Density** of a dynamical system
- 3. Non-trivial Invariant Densities arise in chaotic systems.

On networks:

- 1. On networks, a random walker visits each node with frequency \propto node degree. #groovy
- 2. Equal probability still present: walkers traverse edges with equal frequency. #totallygroovy

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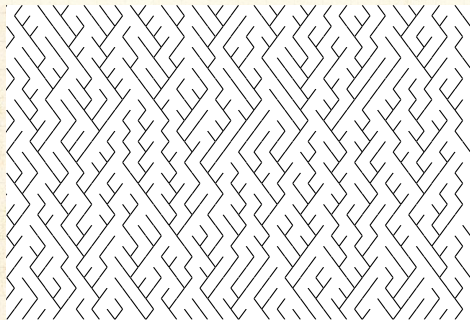
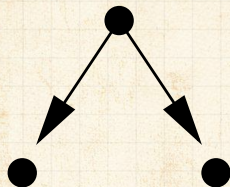
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-  Random directed network on triangular lattice.
-  Toy model of real networks.
-  'Flow' is southeast or southwest with equal probability.

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 Creates basins with random walk boundaries.

 **Observe** that subtracting one random walk from another gives random walk with increments:

$$\epsilon_t = \begin{cases} +1 & \text{with probability } 1/4 \\ 0 & \text{with probability } 1/2 \\ -1 & \text{with probability } 1/4 \end{cases}$$

 Random walk with probabilistic pauses.

 Basin termination = first return random walk problem.

 Basin length l distribution: $P(l) \propto l^{-3/2}$

 For real river networks, generalize to $P(l) \propto l^{-\gamma}$.

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$$\epsilon_t = \begin{cases} +1 & \text{with probability } 1/4 \\ 0 & \text{with probability } 1/2 \\ -1 & \text{with probability } 1/4 \end{cases}$$

 Random walk with probabilistic pauses.

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 For real river networks, generalize to $P(\ell) \propto \ell^{-\gamma}$.

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Connections between exponents:

 For a basin of length l , width $\propto l^{1/2}$

 Basin area $a \propto l \cdot l^{1/2} = l^{3/2}$

 Invert: $l \propto a^{2/3}$

 $dl \propto d(a^{2/3}) = 2/3 a^{-1/3} da$

 $\Pr(\text{basin area} = a) da$
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Connections between exponents:

- Both basin area and length obey power law distributions
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- Reportedly: $1.3 < \tau < 1.5$ and $1.5 < \gamma < 2$

Generalize relationship between area and length.


Hack's law¹⁵¹

$$L \propto A^h$$

- For real, large networks $h \approx 0.5$
- Smaller basins possibly $h > 1/2$ (see: allometry)
- Models exist with interesting values of h .
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$$d\ell \propto d(a^h) = ha^{h-1}da$$

Find τ in terms of γ and h .

$$\Pr(\text{basin area} = a)da \\ = \Pr(\text{basin length} = \ell)d\ell$$

$$\tau = 1 + h(\gamma - 1)$$

Excellent example of the **Scaling Relations** found between exponents describing power laws for many systems.

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Connections between exponents:

With more detailed description of network structure, $\tau = 1 + h(\gamma - 1)$ simplifies to: ^[1]

$$\tau = 2 - h$$

and

$$\gamma = 1/h$$

- Only one exponent is independent (take h).
- Simplifies system description.
- Expect Scaling Relations where power laws are found.
- Need only characterize Universality class with independent exponents.

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Other First Returns or First Passage Times:


PoCS | @pocsvox


Power-Law
Mechanisms, Pt. 1

Failure:

 A very simple model of failure/death: ^[11]

 x_t = entity's 'health' at time t

 Start with $x_0 > 0$.

 Entity fails when x hits 0.

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Streams

 Dispersion of suspended sediments in streams.

 Long times for clearing.





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More than randomness

 Can generalize to Fractional Random Walks [7, 8, 6]

 Levy flights, Fractional Brownian Motion

 See Montroll and Shlesinger for example: [6]
"On $1/f$ noise and other distributions with long tails."

Proc. Natl. Acad. Sci., 1982.

 In 1-d, standard deviation σ scales as

$$\sigma \sim t^\alpha$$

 Extensive memory of path now matters...

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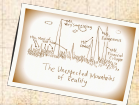
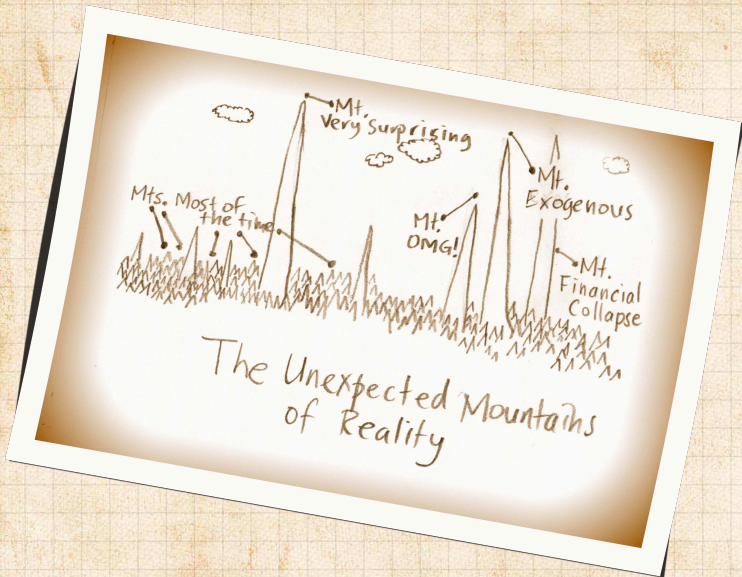
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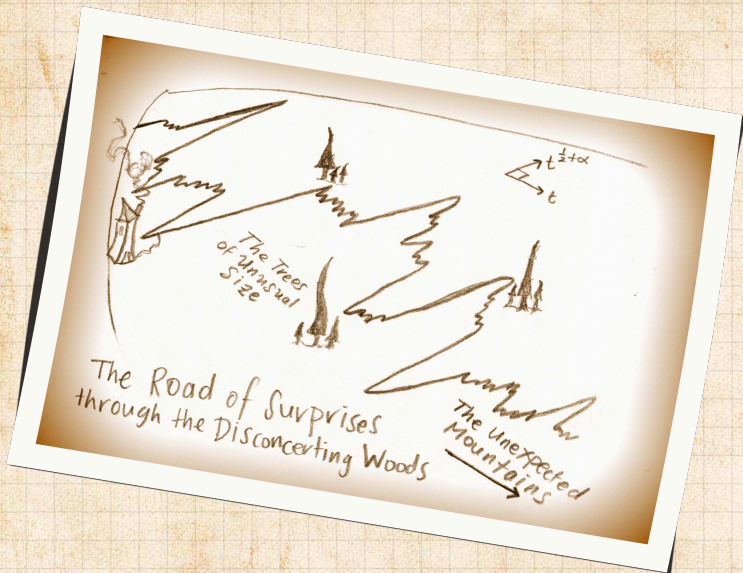
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Neural reboot (NR):

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Desert rain frog/Squeaky toy:

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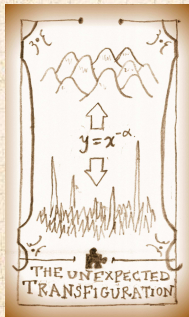
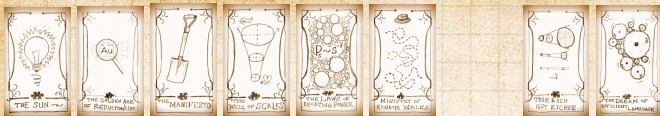
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<https://www.youtube.com/v/cBkWhkAZ9ds?rel=0>





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PoCS | @pocsvox

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Variable Transformation

Understand power laws as arising from

1. Elementary distributions (e.g., exponentials).
2. Variables connected by power relationships.

1. Random variable X with known distribution P_X

2. Second random variable Y with $y = f(x)$.

$$\begin{aligned} P_Y(y)dy &= \\ \sum_{x|f(x)=y} P_X(x)dx &= \\ \sum_{y|f(x)=y} P_X(f^{-1}(y)) \frac{dy}{|f'(f^{-1}(y))|} \end{aligned}$$

3. Often easier to do by hand...

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General Example

1. Assume relationship between x and y is 1-1.

2. Power-law relationship between variables:

$$y = cx^{-\alpha}, \alpha > 0$$

3. Look at y large and x small



$$dy = d(cx^{-\alpha})$$

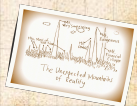
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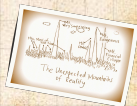
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Now make transformation:

$$P_y(y)dy = P_x(x)dx$$

⚡ If $P_x(x) \rightarrow$ non-zero constant as $x \rightarrow 0$ then

$$P_x(y) \propto y^{-1-1/\alpha} \text{ as } y \rightarrow \infty.$$

⚡ If $P_x(x) \rightarrow x^\beta$ as $x \rightarrow 0$ then

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Now make transformation:

$$P_y(y)dy = P_x(x)dx$$

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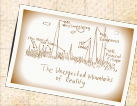
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Example

Exponential distribution

Given $P_x(x) = \frac{1}{\lambda}e^{-x/\lambda}$ and $y = cx^{-\alpha}$, then

$$P(y) \propto y^{-1-1/\alpha} + O(y^{-1-2/\alpha})$$

- Exponentials arise from randomness (easy)...
- More later when we cover robustness.

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Gravity

PoCS | @pocsvox

Power-Law
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☰ Select a random point in the universe \vec{x}

☞ Measure the force of gravity $F(\vec{x})$

☞ Observe that $P_F(F) \sim F^{-5/2}$.



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Gravity

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 **PoCS**
Principles of
Complex Systems
@pocsvox
What's the Story?



Matter is concentrated in stars: [10]

 F is distributed unevenly

 Probability of being a distance r from a single star at $\vec{x} = \vec{0}$:

$$P_r(r)dr \propto r^2 dr$$

 Assume stars are distributed randomly in space (oops?)

 Assume only one star has significant effect at \vec{x} .

 Law of gravity:

$$F \propto r^{-2}$$

 invert:

$$r \propto F^{-\frac{1}{2}}$$

 Connect differentials: $dr \propto dF^{-\frac{1}{2}} \propto F^{-\frac{3}{2}} dF$


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Transformation:

Using $r \propto F^{-1/2}$, $dr \propto F^{-3/2}dF$, and $P_r(r) \propto r^2$

$$P_F(F)dF = P_r(r)dr$$

$$\propto P_r(\text{const} \times F^{-1/2})F^{-3/2}dF$$

$$\propto (F^{-1/2})^2 F^{-3/2}dF$$

$$= F^{-1-3/2}dF$$

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Gravity:

$$P_F(F) = F^{-5/2} dF$$

$$\gamma = 5/2$$

- Mean is finite.
- Variance = ∞ .
- A wild distribution.
- Upshot: Random sampling of space usually safe but can end badly...

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Variance = ∞ .



A wild distribution.



Upshot: Random sampling of space usually safe
but can end badly...

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Doctorin' the Tardis

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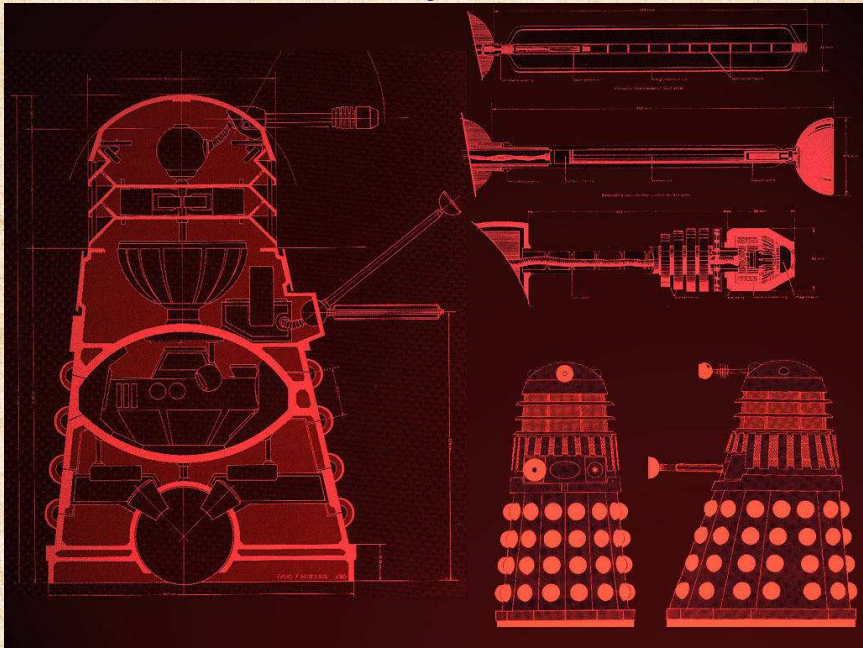
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□ Todo: Build Dalek army.



Outline

PoCS | @pocsvox

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Mechanisms, Pt. 1

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PLIPLO = **Power law in, power law out**



Explain a power law as resulting from another unexplained power law.



Yet another homunculus argument?...



Don't do this!!! (slap, slap)



MIWO = **Mild in, Wild out** is the stuff.



In general: We need mechanisms!



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Neural reboot (NR):

PoCS | @pocsvox

Power-Law
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Zoomage in slow motion

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